

A partial Laplacian as an infinitesimal generator on the Wasserstein space

YT. Chow and W. Gangbo
NIPS 17 - OTML

December 9, 2017

- ▶ Let $\mathcal{P}_2(\mathbb{R}^d)$ be the set of probability measures on \mathbb{R}^d , of finite second moments

- ▶ Let $\mathcal{P}_2(\mathbb{R}^d)$ be the set of probability measures on \mathbb{R}^d , of finite second moments
- ▶ Goal 1:

- ▶ Let $\mathcal{P}_2(\mathbb{R}^d)$ be the set of probability measures on \mathbb{R}^d , of finite second moments
- ▶ Goal 1:

Define Brownian motions on $\mathcal{P}_2(\mathbb{R}^d)$



▶ Let $\mathcal{P}_2(\mathbb{R}^d)$ be the set of probability measures on \mathbb{R}^d , of finite second moments

▶ Goal 1:

Define Brownian motions on $\mathcal{P}_2(\mathbb{R}^d)$

▶

▶

▶ Let $\mathcal{P}_2(\mathbb{R}^d)$ be the set of probability measures on \mathbb{R}^d , of finite second moments

▶ Goal 1:

Define Brownian motions on $\mathcal{P}_2(\mathbb{R}^d)$

▶

▶

▶ Goal 2:

▶ Let $\mathcal{P}_2(\mathbb{R}^d)$ be the set of probability measures on \mathbb{R}^d , of finite second moments

▶ Goal 1:

Define Brownian motions on $\mathcal{P}_2(\mathbb{R}^d)$

▶

▶

▶ Goal 2:

Convince you we have the right definition

▶

▶ Let $\mathcal{P}_2(\mathbb{R}^d)$ be the set of probability measures on \mathbb{R}^d , of finite second moments

▶ Goal 1:

Define Brownian motions on $\mathcal{P}_2(\mathbb{R}^d)$

▶

▶

▶ Goal 2:

Convince you we have the right definition

▶

▶

▶ Let $\mathcal{P}_2(\mathbb{R}^d)$ be the set of probability measures on \mathbb{R}^d , of finite second moments

▶ Goal 1:

Define Brownian motions on $\mathcal{P}_2(\mathbb{R}^d)$

▶

▶

▶ Goal 2:

Convince you we have the right definition

▶

▶

▶ Goal 3:

▶ Let $\mathcal{P}_2(\mathbb{R}^d)$ be the set of probability measures on \mathbb{R}^d , of finite second moments

▶ Goal 1:

Define Brownian motions on $\mathcal{P}_2(\mathbb{R}^d)$

▶

▶

▶ Goal 2:

Convince you we have the right definition

▶

▶

▶ Goal 3:

Convince you of their useful

▶

Goal 3: Example

- ▶ For each integer $n \geq 1$ assume we solve

$$\partial_t U^n(t, x) = \sum_{j=1}^n \Delta_{x_j} U^n(t, x), \quad t > 0, \quad x \in \mathbb{R}^d \times \cdots \times \mathbb{R}^d \quad (1)$$

Goal 3: Example

- ▶ For each integer $n \geq 1$ assume we solve

$$\partial_t U^n(t, x) = \sum_{j=1}^n \Delta_{x_j} U^n(t, x), \quad t > 0, \quad x \in \mathbb{R}^d \times \cdots \times \mathbb{R}^d \quad (1)$$

- ▶ Assume $U^n(t, \cdot)$ is invariant under permutation of the x_i :

$$U^n(t, x_1, \dots, x_{n-1}, x_n) = U^n(t, x_1, \dots, x_n, x_{n-1})$$

Goal 3: Example

- ▶ For each integer $n \geq 1$ assume we solve

$$\partial_t U^n(t, x) = \sum_{j=1}^n \Delta_{x_j} U^n(t, x), \quad t > 0, \quad x \in \mathbb{R}^d \times \cdots \times \mathbb{R}^d \quad (1)$$

- ▶ Assume $U^n(t, \cdot)$ is invariant under permutation of the x_i :

$$U^n(t, x_1, \dots, x_{n-1}, x_n) = U^n(t, x_1, \dots, x_n, x_{n-1})$$

- ▶ We want to make sense of

$$U^n(t, \cdot) \rightarrow U, \quad \text{as } n \rightarrow \infty$$

Goal 3: Example

- ▶ For each integer $n \geq 1$ assume we solve

$$\partial_t U^n(t, x) = \sum_{j=1}^n \Delta_{x_j} U^n(t, x), \quad t > 0, \quad x \in \mathbb{R}^d \times \cdots \times \mathbb{R}^d \quad (1)$$

- ▶ Assume $U^n(t, \cdot)$ is invariant under permutation of the x_i :

$$U^n(t, x_1, \dots, x_{n-1}, x_n) = U^n(t, x_1, \dots, x_n, x_{n-1})$$

- ▶ We want to make sense of

$$U^n(t, \cdot) \rightarrow U, \quad \text{as } n \rightarrow \infty$$

- ▶ and check what equation U satisfies:

$$\partial_t U = \sum_{j=1}^{\infty} \Delta_{x_j} U?$$

Review: geometry and stochastic

- ▶ Let \mathbb{M} be a finite dimensional euclidean space (e.g. $\mathbb{M} = \mathbb{R}^d$).

Review: geometry and stochastic

- ▶ Let \mathbb{M} be a finite dimensional euclidean space (e.g. $\mathbb{M} = \mathbb{R}^d$).
- ▶ For $x \in \mathbb{M}$, let $T_x\mathbb{M}$ denote the tangent space at x

Review: geometry and stochastic

- ▶ Let \mathbb{M} be a finite dimensional euclidean space (e.g. $\mathbb{M} = \mathbb{R}^d$).
- ▶ For $x \in \mathbb{M}$, let $T_x\mathbb{M}$ denote the tangent space at x
- ▶ The matrix of the second derivatives of $g \in C^2(\mathbb{M})$ define a bilinear form

$$\text{hess } g(x) : T_x\mathbb{M} \times T_x\mathbb{M} \rightarrow \mathbb{R} :$$

$$\text{hess } g(x)(v, v) = \left. \frac{d^2}{dt^2} g(x + tv) \right|_{t=0}$$

Review: geometry and stochastic

- ▶ Let \mathbb{M} be a finite dimensional euclidean space (e.g. $\mathbb{M} = \mathbb{R}^d$).
- ▶ For $x \in \mathbb{M}$, let $T_x\mathbb{M}$ denote the tangent space at x
- ▶ The matrix of the second derivatives of $g \in C^2(\mathbb{M})$ define a bilinear form

$$\text{hess } g(x) : T_x\mathbb{M} \times T_x\mathbb{M} \rightarrow \mathbb{R} :$$

$$\text{hess } g(x)(v, v) = \left. \frac{d^2}{dt^2} g(x + tv) \right|_{t=0}$$

- ▶ It has been obtained by the 2nd derivatives of g along the geodesic $t \rightarrow x + tv$

Review: geometry and stochastic

- ▶ Let \mathbb{M} be a finite dimensional euclidean space (e.g. $\mathbb{M} = \mathbb{R}^d$).
- ▶ For $x \in \mathbb{M}$, let $T_x\mathbb{M}$ denote the tangent space at x
- ▶ The matrix of the second derivatives of $g \in C^2(\mathbb{M})$ define a bilinear form

$$\text{hess } g(x) : T_x\mathbb{M} \times T_x\mathbb{M} \rightarrow \mathbb{R} :$$

$$\text{hess } g(x)(v, v) = \left. \frac{d^2}{dt^2} g(x + tv) \right|_{t=0}$$

- ▶ It has been obtained by the 2nd derivatives of g along the geodesic $t \rightarrow x + tv$
- ▶ We have

$$\Delta g(x) = \sum_{j=1}^d \text{hess } g(x)(v_j, v_j)$$

Review: geometry and stochastic

- ▶ Let \mathbb{M} be a finite dimensional euclidean space (e.g. $\mathbb{M} = \mathbb{R}^d$).
- ▶ For $x \in \mathbb{M}$, let $T_x\mathbb{M}$ denote the tangent space at x
- ▶ The matrix of the second derivatives of $g \in C^2(\mathbb{M})$ define a bilinear form

$$\text{hess } g(x) : T_x\mathbb{M} \times T_x\mathbb{M} \rightarrow \mathbb{R} :$$

$$\text{hess } g(x)(v, v) = \left. \frac{d^2}{dt^2} g(x + tv) \right|_{t=0}$$

- ▶ It has been obtained by the 2nd derivatives of g along the geodesic $t \rightarrow x + tv$
- ▶ We have

$$\Delta g(x) = \sum_{j=1}^d \text{hess } g(x)(v_j, v_j)$$

- ▶ Here $\{v_j\}_{j=1}^d$ is any orthonormal basis of $T_x\mathbb{M}$

Review: geometry and stochastic

- ▶ Let $(W_t)_{t \geq 0}$ be a Brownian motion on \mathbb{R}^d

Review: geometry and stochastic

- ▶ Let $(W_t)_{t \geq 0}$ be a Brownian motion on \mathbb{R}^d
- ▶ We have

$$\Delta g(x) = \lim_{t \rightarrow 0} \frac{\mathbb{E}(g(x + \sqrt{2}W_t)) - g(x)}{t}$$

Review: geometry and stochastic

- ▶ Let $(W_t)_{t \geq 0}$ be a Brownian motion on \mathbb{R}^d
- ▶ We have

$$\Delta g(x) = \lim_{t \rightarrow 0} \frac{\mathbb{E}(g(x + \sqrt{2}W_t)) - g(x)}{t}$$

- ▶ This is related to Feynman-Kac formula: if we set

$$u(t, x) = \mathbb{E}(g(x + \sqrt{2}W_t))$$

Review: geometry and stochastic

- ▶ Let $(W_t)_{t \geq 0}$ be a Brownian motion on \mathbb{R}^d
- ▶ We have

$$\Delta g(x) = \lim_{t \rightarrow 0} \frac{\mathbb{E}(g(x + \sqrt{2}W_t)) - g(x)}{t}$$

- ▶ This is related to Feynman-Kac formula: if we set

$$u(t, x) = \mathbb{E}(g(x + \sqrt{2}W_t))$$

- ▶ then

$$\partial_t u = \Delta u, \quad u(0, \cdot) = g$$

EXCURSION TO WASSERSTEIN GEOMETRY

- ▶ Our first goal is to extend the above considerations to the Wasserstein space $\mathcal{P}_2(\mathbb{R}^d)$.

EXCURSION TO WASSERSTEIN GEOMETRY

- ▶ Our first goal is to extend the above considerations to the Wasserstein space $\mathcal{P}_2(\mathbb{R}^d)$.
- ▶ This includes computing Hessians on $\mathcal{P}_2(\mathbb{R}^d)$

EXCURSION TO WASSERSTEIN GEOMETRY

- ▶ Our first goal is to extend the above considerations to the Wasserstein space $\mathcal{P}_2(\mathbb{R}^d)$.
- ▶ This includes computing Hessians on $\mathcal{P}_2(\mathbb{R}^d)$
- ▶ The first annoying complication to describe Hessians is

EXCURSION TO WASSERSTEIN GEOMETRY

- ▶ Our first goal is to extend the above considerations to the Wasserstein space $\mathcal{P}_2(\mathbb{R}^d)$.
- ▶ This includes computing Hessians on $\mathcal{P}_2(\mathbb{R}^d)$
- ▶ The first annoying complication to describe Hessians is
- ▶ $\mathcal{P}_2(\mathbb{R}^d)$ is not a euclidean space

EXCURSION TO WASSERSTEIN GEOMETRY

- ▶ Our first goal is to extend the above considerations to the Wasserstein space $\mathcal{P}_2(\mathbb{R}^d)$.
- ▶ This includes computing Hessians on $\mathcal{P}_2(\mathbb{R}^d)$
- ▶ The first annoying complication to describe Hessians is
- ▶ $\mathcal{P}_2(\mathbb{R}^d)$ is not a euclidean space
- ▶ Life would be simpler on a Hilbert space such as

$$\mathbb{H} := L^2(\Omega, \mathbb{R}^d), \quad \Omega := (0, 1)^d \subset \mathbb{R}^d.$$

EXCURSION TO WASSERSTEIN GEOMETRY

- ▶ Our first goal is to extend the above considerations to the Wasserstein space $\mathcal{P}_2(\mathbb{R}^d)$.
- ▶ This includes computing Hessians on $\mathcal{P}_2(\mathbb{R}^d)$
- ▶ The first annoying complication to describe Hessians is
- ▶ $\mathcal{P}_2(\mathbb{R}^d)$ is not a Euclidean space
- ▶ Life would be simpler on a Hilbert space such as

$$\mathbb{H} := L^2(\Omega, \mathbb{R}^d), \quad \Omega := (0, 1)^d \subset \mathbb{R}^d.$$

- ▶ for future studies, we showed (G-Tudorascu):

EXCURSION TO WASSERSTEIN GEOMETRY

- ▶ Our first goal is to extend the above considerations to the Wasserstein space $\mathcal{P}_2(\mathbb{R}^d)$.
- ▶ This includes computing Hessians on $\mathcal{P}_2(\mathbb{R}^d)$
- ▶ The first annoying complication to describe Hessians is
- ▶ $\mathcal{P}_2(\mathbb{R}^d)$ is not a Euclidean space
- ▶ Life would be simpler on a Hilbert space such as

$$\mathbb{H} := L^2(\Omega, \mathbb{R}^d), \quad \Omega := (0, 1)^d \subset \mathbb{R}^d.$$

- ▶ for future studies, we showed (G-Tudorascu):
- ▶ the differential structure on $\mathcal{P}_2(\mathbb{R}^d)$, by Ambrosio–Gigli–Savaré

EXCURSION TO WASSERSTEIN GEOMETRY

- ▶ Our first goal is to extend the above considerations to the Wasserstein space $\mathcal{P}_2(\mathbb{R}^d)$.
- ▶ This includes computing Hessians on $\mathcal{P}_2(\mathbb{R}^d)$
- ▶ The first annoying complication to describe Hessians is
- ▶ $\mathcal{P}_2(\mathbb{R}^d)$ is not a Euclidean space
- ▶ Life would be simpler on a Hilbert space such as

$$\mathbb{H} := L^2(\Omega, \mathbb{R}^d), \quad \Omega := (0, 1)^d \subset \mathbb{R}^d.$$

- ▶ for future studies, we showed (G-Tudorascu):
- ▶ the differential structure on $\mathcal{P}_2(\mathbb{R}^d)$, by Ambrosio–Gigli–Savaré
- ▶ **amount**

$$\mathcal{P}_2(\mathbb{R}^d) = \mathbb{H} / \text{equi relation}$$

▶ Lift $G : \mathcal{P}_2(\mathbb{R}^d) \rightarrow \mathbb{R}$ to obtain $\bar{G} : \mathbb{H} \rightarrow \mathbb{R}$

▶ Lift $G : \mathcal{P}_2(\mathbb{R}^d) \rightarrow \mathbb{R}$ to obtain $\bar{G} : \mathbb{H} \rightarrow \mathbb{R}$



$$\text{"Hess}_{\mathcal{P}_2(\mathbb{R}^d)} G = \text{Hess}_{\mathbb{H}} \bar{G}\text{"}$$

▶ Lift $G : \mathcal{P}_2(\mathbb{R}^d) \rightarrow \mathbb{R}$ to obtain $\bar{G} : \mathbb{H} \rightarrow \mathbb{R}$

▶

$$\text{"Hess}_{\mathcal{P}_2(\mathbb{R}^d)} G = \text{Hess}_{\mathbb{H}} \bar{G}\text{"}$$

▶ Recall: if

$$\mathbf{v} := \nabla \varphi \in \nabla C_c^\infty\left((-1, 1) \times \mathbb{R}^d, \mathbb{R}^d\right)$$

- ▶ Lift $G : \mathcal{P}_2(\mathbb{R}^d) \rightarrow \mathbb{R}$ to obtain $\bar{G} : \mathbb{H} \rightarrow \mathbb{R}$



$$\text{"Hess}_{\mathcal{P}_2(\mathbb{R}^d)} G = \text{Hess}_{\mathbb{H}} \bar{G}\text{"}$$

- ▶ Recall: if

$$\mathbf{v} := \nabla \varphi \in \nabla C_c^\infty\left((-1, 1) \times \mathbb{R}^d, \mathbb{R}^d\right)$$

- ▶ given $m \in \mathcal{P}_2(\mathbb{R}^d)$, we have a geodesic on $\mathcal{P}_2(\mathbb{R}^d)$ starting at m :

$$\partial_t \sigma + \nabla \cdot (\sigma \mathbf{v}) = 0, \quad \sigma(0) = m$$

- ▶ Lift $G : \mathcal{P}_2(\mathbb{R}^d) \rightarrow \mathbb{R}$ to obtain $\bar{G} : \mathbb{H} \rightarrow \mathbb{R}$



$$\text{"Hess}_{\mathcal{P}_2(\mathbb{R}^d)} G = \text{Hess}_{\mathbb{H}} \bar{G}\text{"}$$

- ▶ Recall: if

$$\mathbf{v} := \nabla \varphi \in \nabla C_c^\infty\left((-1, 1) \times \mathbb{R}^d, \mathbb{R}^d\right)$$

- ▶ given $m \in \mathcal{P}_2(\mathbb{R}^d)$, we have a geodesic on $\mathcal{P}_2(\mathbb{R}^d)$ starting at m :

$$\partial_t \sigma + \nabla \cdot (\sigma \mathbf{v}) = 0, \quad \sigma(0) = m$$

- ▶ The Hessian of $G : \mathcal{P}_2(\mathbb{R}^d) \rightarrow \mathbb{R}$ at m is a bilinear form

- ▶ Lift $G : \mathcal{P}_2(\mathbb{R}^d) \rightarrow \mathbb{R}$ to obtain $\bar{G} : \mathbb{H} \rightarrow \mathbb{R}$



$$\text{"Hess}_{\mathcal{P}_2(\mathbb{R}^d)} G = \text{Hess}_{\mathbb{H}} \bar{G}\text{"}$$

- ▶ Recall: if

$$\mathbf{v} := \nabla\varphi \in \nabla C_c^\infty\left((-1, 1) \times \mathbb{R}^d, \mathbb{R}^d\right)$$

- ▶ given $m \in \mathcal{P}_2(\mathbb{R}^d)$, we have a geodesic on $\mathcal{P}_2(\mathbb{R}^d)$ starting at m :

$$\partial_t \sigma + \nabla \cdot (\sigma \mathbf{v}) = 0, \quad \sigma(0) = m$$

- ▶ The Hessian of $G : \mathcal{P}_2(\mathbb{R}^d) \rightarrow \mathbb{R}$ at m is a bilinear form



$$\text{Hess} G[m](\xi, \xi) = \left. \frac{d^2}{dt^2} G(\sigma(t)) \right|_{t=0}, \quad \xi := \mathbf{v}(0, \cdot)$$

- ▶ Lift $G : \mathcal{P}_2(\mathbb{R}^d) \rightarrow \mathbb{R}$ to obtain $\bar{G} : \mathbb{H} \rightarrow \mathbb{R}$



$$\text{"Hess}_{\mathcal{P}_2(\mathbb{R}^d)} G = \text{Hess}_{\mathbb{H}} \bar{G}\text{"}$$

- ▶ Recall: if

$$\mathbf{v} := \nabla\varphi \in \nabla C_c^\infty\left((-1, 1) \times \mathbb{R}^d, \mathbb{R}^d\right)$$

- ▶ given $m \in \mathcal{P}_2(\mathbb{R}^d)$, we have a geodesic on $\mathcal{P}_2(\mathbb{R}^d)$ starting at m :

$$\partial_t \sigma + \nabla \cdot (\sigma \mathbf{v}) = 0, \quad \sigma(0) = m$$

- ▶ The Hessian of $G : \mathcal{P}_2(\mathbb{R}^d) \rightarrow \mathbb{R}$ at m is a bilinear form



$$\text{Hess} G[m](\xi, \xi) = \left. \frac{d^2}{dt^2} G(\sigma(t)) \right|_{t=0}, \quad \xi := \mathbf{v}(0, \cdot)$$



EXCURSION IN STOCHASTIC WORLD

- ▶ For a while, stop with the above deterministic considerations

EXCURSION IN STOCHASTIC WORLD

- ▶ For a while, stop with the above deterministic considerations
- ▶ Choose a Brownian motion $(W_t)_{t \geq 0}$ on \mathbb{R}^d

EXCURSION IN STOCHASTIC WORLD

- ▶ For a while, stop with the above deterministic considerations
- ▶ Choose a Brownian motion $(W_t)_{t \geq 0}$ on \mathbb{R}^d
- ▶ Lift it to $\mathcal{P}_2(\mathbb{R}^d)$ as follows

EXCURSION IN STOCHASTIC WORLD

- ▶ For a while, stop with the above deterministic considerations
- ▶ Choose a Brownian motion $(W_t)_{t \geq 0}$ on \mathbb{R}^d
- ▶ Lift it to $\mathcal{P}_2(\mathbb{R}^d)$ as follows



$$t \rightarrow \mathbb{B}_t^m := \left(\mathbf{id} + \sqrt{2}W_t \right) \# m$$

EXCURSION IN STOCHASTIC WORLD

- ▶ For a while, stop with the above deterministic considerations
- ▶ Choose a Brownian motion $(W_t)_{t \geq 0}$ on \mathbb{R}^d
- ▶ Lift it to $\mathcal{P}_2(\mathbb{R}^d)$ as follows



$$t \rightarrow \mathbb{B}_t^m := \left(\mathbf{id} + \sqrt{2}W_t \right) \# m$$

- ▶ Here \mathbf{id} is the identity map and so, \mathbb{B}^m is a stochastic motion in $\mathcal{P}_2(\mathbb{R}^d)$, starting at m

EXCURSION IN STOCHASTIC WORLD

- ▶ For a while, stop with the above deterministic considerations
- ▶ Choose a Brownian motion $(W_t)_{t \geq 0}$ on \mathbb{R}^d
- ▶ Lift it to $\mathcal{P}_2(\mathbb{R}^d)$ as follows



$$t \rightarrow \mathbb{B}_t^m := \left(\mathbf{id} + \sqrt{2}W_t \right) \# m$$

- ▶ Here \mathbf{id} is the identity map and so, \mathbb{B}^m is a stochastic motion in $\mathcal{P}_2(\mathbb{R}^d)$, starting at m



EXCURSION IN STOCHASTIC WORLD

- ▶ For a while, stop with the above deterministic considerations
- ▶ Choose a Brownian motion $(W_t)_{t \geq 0}$ on \mathbb{R}^d
- ▶ Lift it to $\mathcal{P}_2(\mathbb{R}^d)$ as follows



$$t \rightarrow \mathbb{B}_t^m := \left(\mathbf{id} + \sqrt{2}W_t \right) \# m$$

- ▶ Here \mathbf{id} is the identity map and so, \mathbb{B}^m is a stochastic motion in $\mathcal{P}_2(\mathbb{R}^d)$, starting at m



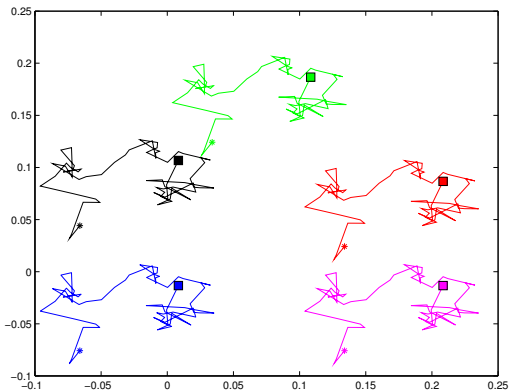
- ▶ For example

$$\mathbb{B}_t^m = \frac{1}{5} \sum_{i=1}^5 \delta_{x_i + \sqrt{2}W_t}$$

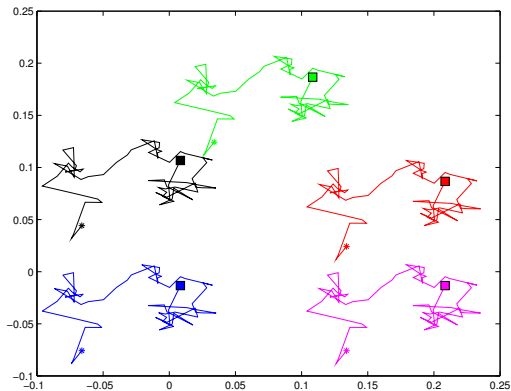
is a population of 5 players evolving with a common noise

- ▶ The Brownian motion is represented in the figure

- ▶ The Brownian motion is represented in the figure

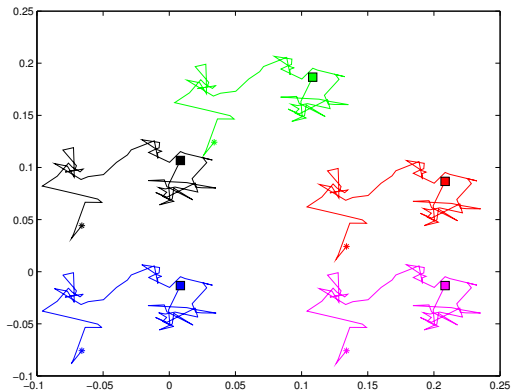


- ▶ The Brownian motion is represented in the figure



- ▶ Notice all the point masses perform a coherent/identical Brownian motion.

- ▶ The Brownian motion is represented in the figure



- ▶ Notice all the point masses perform a coherent/identical Brownian motion.



Linking geometry to stochastic

- ▶ Let $\{E_j\}_{j=1}^d$ be an orthonormal basis in \mathbb{R}^d . We define the partial Wasserstein Laplacian at m as the trace:

$$\Delta_w G[m] = \sum_{j=1}^d \text{Hess } G[m](E_j, E_j)$$

Linking geometry to stochastic

- ▶ Let $\{E_j\}_{j=1}^d$ be an orthonormal basis in \mathbb{R}^d . We define the partial Wasserstein Laplacian at m as the trace:

$$\Delta_w G[m] = \sum_{j=1}^d \text{Hess } G[m](E_j, E_j)$$

- ▶ The first argument supporting the fact that we have the right definition is

$$\Delta_w G[m] = \lim_{t \rightarrow 0} \frac{\mathbb{E}(G(\mathbb{B}_t^m)) - G(m)}{t}$$

Linking geometry to stochastic

- ▶ Let $\{E_j\}_{j=1}^d$ be an orthonormal basis in \mathbb{R}^d . We define the partial Wasserstein Laplacian at m as the trace:

$$\Delta_w G[m] = \sum_{j=1}^d \text{Hess } G[m](E_j, E_j)$$

- ▶ The first argument supporting the fact that we have the right definition is

$$\Delta_w G[m] = \lim_{t \rightarrow 0} \frac{\mathbb{E}(G(\mathbb{B}_t^m)) - G(m)}{t}$$

- ▶ Another argument related to Feynman–Kac: set

$$U(t, m) := E(G(\mathbb{B}_t^m))$$

to discover that

Linking geometry to stochastic

- ▶ Let $\{E_j\}_{j=1}^d$ be an orthonormal basis in \mathbb{R}^d . We define the partial Wasserstein Laplacian at m as the trace:

$$\Delta_w G[m] = \sum_{j=1}^d \text{Hess } G[m](E_j, E_j)$$

- ▶ The first argument supporting the fact that we have the right definition is

$$\Delta_w G[m] = \lim_{t \rightarrow 0} \frac{\mathbb{E}(G(\mathbb{B}_t^m)) - G(m)}{t}$$

- ▶ Another argument related to Feynman–Kac: set

$$U(t, m) := E(G(\mathbb{B}_t^m))$$

to discover that



$$\partial_t U = \Delta_w U, \quad U(0, \cdot) = G$$

Example: system of n particles

- ▶ Use the equation

$$\partial_t U = \Delta_w U, \quad U(0, \cdot) = G$$

Example: system of n particles

- ▶ Use the equation

$$\partial_t U = \Delta_w U, \quad U(0, \cdot) = G$$

- ▶ Set

$$G^n(x_1, \dots, x_n) = G\left(\frac{1}{n} \sum_{j=1}^n \delta_{x_j}\right).$$

and

$$U\left(t, \frac{1}{n} \sum_{j=1}^n \delta_{x_j}\right) = U^n(t, x_1, \dots, x_n)$$

Example: system of n particles

- ▶ Use the equation

$$\partial_t U = \Delta_w U, \quad U(0, \cdot) = G$$

- ▶ Set

$$G^n(x_1, \dots, x_n) = G\left(\frac{1}{n} \sum_{j=1}^n \delta_{x_j}\right).$$

and

$$U\left(t, \frac{1}{n} \sum_{j=1}^n \delta_{x_j}\right) = U^n(t, x_1, \dots, x_n)$$

- ▶ Then U^n satisfies the PDE

$$\partial_t U^n = \sum_{j,k=1}^n \text{Trace} \left(\nabla_{x_j x_k} U^n \right), \quad U^n(0, \cdot) = G^n$$

$n \rightarrow \infty$

- ▶ Observe the decomposition

$$\sum_{j,k=1}^n \text{Trace} \left(\nabla_{x_j x_k} G^n \right) = \sum_j^n \Delta_{x_j} G^n + \sum_{j \neq k}^n \text{Trace} \left(\nabla_{x_j x_k} G^n \right)$$

- ▶ Observe the decomposition

$$\sum_{j,k=1}^n \text{Trace} \left(\nabla_{x_j x_k} G^n \right) = \sum_j^n \Delta_{x_j} G^n + \sum_{j \neq k}^n \text{Trace} \left(\nabla_{x_j x_k} G^n \right)$$

- ▶ Let $\nabla_w G$ and $\nabla_{ww} G$ be the first and second order Wasserstein gradients.

$n \rightarrow \infty$

- ▶ Observe the decomposition

$$\sum_{j,k=1}^n \text{Trace} \left(\nabla_{x_j x_k} G^n \right) = \sum_j \Delta_{x_j} G^n + \sum_{j \neq k} \text{Trace} \left(\nabla_{x_j x_k} G^n \right)$$

- ▶ Let $\nabla_w G$ and $\nabla_{ww} G$ be the first and second order Wasserstein gradients.
- ▶ We have

$$\sum_{j=1}^n \Delta_{x_j} G^n \rightarrow \int_{\mathbb{R}^d} \text{div}_x (\nabla_w G[m](x)) m(dx).$$

if

$$\frac{1}{n} \sum_{j=1}^n \delta_{x_j} \rightarrow m$$

- ▶ Observe the decomposition

$$\sum_{j,k=1}^n \text{Trace} \left(\nabla_{x_j x_k} G^n \right) = \sum_j \Delta_{x_j} G^n + \sum_{j \neq k} \text{Trace} \left(\nabla_{x_j x_k} G^n \right)$$

- ▶ Let $\nabla_w G$ and $\nabla_{ww} G$ be the first and second order Wasserstein gradients.
- ▶ We have

$$\sum_{j=1}^n \Delta_{x_j} G^n \rightarrow \int_{\mathbb{R}^d} \text{div}_x (\nabla_w G[m](x)) m(dx).$$

if

$$\frac{1}{n} \sum_{j=1}^n \delta_{x_j} \rightarrow m$$

- ▶ We have

$$\sum_{j \neq k} \text{Trace} \left(\nabla_{x_j x_k} G^n \right) \rightarrow \int_{\mathbb{R}^{2d}} \text{Tr} \left(\nabla_w^2 G[m](x, a) \right) m(dx) m(da)$$

Partial Laplacian

- ▶ Decompose Δ_w into the sum of two operators:

Partial Laplacian

- ▶ Decompose Δ_w into the sum of two operators:



$$\Delta_w G[m] = O(G)(m) + \int_{\mathbb{R}^{2d}} \text{Tr} \left(\nabla_w^2 G[m](x, a) \right) m(dx) m(da)$$

where

$$O(G)(m) := \int_{\mathbb{R}^d} \text{div}_x (\nabla_w G[m](x)) m(dx).$$

Partial Laplacian

- ▶ Decompose Δ_w into the sum of two operators:



$$\Delta_w G[m] = O(G)(m) + \int_{\mathbb{R}^{2d}} \text{Tr} \left(\nabla_w^2 G[m](x, a) \right) m(dx) m(da)$$

where

$$O(G)(m) := \int_{\mathbb{R}^d} \text{div}_x (\nabla_w G[m](x)) m(dx).$$

- ▶ Eigenvalues of $-\Delta_w$ in Fourier

$$\lambda_k^2 := 4\pi^2 \left\| \sum_{j=1}^d \xi_j \right\|^2 = 4\pi^2 \sum_{j=1}^d \|\xi_j\|^2 - 4\pi^2 \sum_{j \neq k} \xi_j \cdot \xi_k$$

Partial Laplacian

- ▶ Decompose Δ_w into the sum of two operators:



$$\Delta_w G[m] = O(G)(m) + \int_{\mathbb{R}^{2d}} \text{Tr} \left(\nabla_w^2 G[m](x, a) \right) m(dx) m(da)$$

where

$$O(G)(m) := \int_{\mathbb{R}^d} \text{div}_x (\nabla_w G[m](x)) m(dx).$$

- ▶ Eigenvalues of $-\Delta_w$ in Fourier

$$\lambda_k^2 := 4\pi^2 \left\| \sum_{j=1}^d \xi_j \right\|^2 = 4\pi^2 \sum_{j=1}^d \|\xi_j\|^2 - 4\pi^2 \sum_{j \neq k} \xi_j \cdot \xi_k$$

- ▶ The eigenvalues of O :

$$-4\pi^2 \sum_{j=1}^d \|\xi_j\|^2$$

and so, $-O$ is a positive operator responsible for smoothing effects

Partial Laplacian

- ▶ There is a perturbation of Δ_w which is uniformly elliptic:

Partial Laplacian

- ▶ There is a perturbation of Δ_w which is uniformly elliptic:
- ▶ Set

$$\Delta_{w,\epsilon} G[m] = (1 + \epsilon)O(G)(m) + \int_{\mathbb{R}^{2d}} \text{Tr}\left(\nabla_w^2 G[m](x, a)\right) m(dx)m(da)$$

Partial Laplacian

- ▶ There is a perturbation of Δ_w which is uniformly elliptic:
- ▶ Set

$$\Delta_{w,\epsilon} G[m] = (1 + \epsilon)O(G)(m) + \int_{\mathbb{R}^{2d}} \text{Tr}\left(\nabla_w^2 G[m](x, a)\right) m(dx)m(da)$$

- ▶ In other words

$$\Delta_{w,\epsilon} = \Delta_w + \epsilon O$$

Partial Laplacian

- ▶ There is a perturbation of Δ_w which is uniformly elliptic:
- ▶ Set

$$\Delta_{w,\epsilon} G[m] = (1 + \epsilon)O(G)(m) + \int_{\mathbb{R}^{2d}} \text{Tr}\left(\nabla_w^2 G[m](x, a)\right) m(dx)m(da)$$

- ▶ In other words

$$\Delta_{w,\epsilon} = \Delta_w + \epsilon O$$

- ▶ $\Delta_{w,\epsilon}$ has a better smoothing effect

Partial Laplacian

- ▶ There is a perturbation of Δ_w which is uniformly elliptic:
- ▶ Set

$$\Delta_{w,\epsilon} G[m] = (1 + \epsilon)O(G)(m) + \int_{\mathbb{R}^{2d}} \text{Tr}\left(\nabla_w^2 G[m](x, a)\right) m(dx)m(da)$$

- ▶ In other words

$$\Delta_{w,\epsilon} = \Delta_w + \epsilon O$$

- ▶ $\Delta_{w,\epsilon}$ has a better smoothing effect
- ▶ Eigenvalues of $-\Delta_{w,\epsilon}$ in Fourier are

$$\lambda_{k,\epsilon}^2 := \epsilon 4\pi^2 \sum_{j=1}^d \|\xi_j\|^2 + 4\pi^2 \left\| \sum_{j=1}^d \xi_j \right\|^2 \geq \epsilon 4\pi^2 \sum_{j=1}^d \|\xi_j\|^2$$

Partial Laplacian

- ▶ There is a perturbation of Δ_w which is uniformly elliptic:
- ▶ Set

$$\Delta_{w,\epsilon} G[m] = (1 + \epsilon)O(G)(m) + \int_{\mathbb{R}^{2d}} \text{Tr}\left(\nabla_w^2 G[m](x, a)\right) m(dx) m(da)$$

- ▶ In other words

$$\Delta_{w,\epsilon} = \Delta_w + \epsilon O$$

- ▶ $\Delta_{w,\epsilon}$ has a better smoothing effect
- ▶ Eigenvalues of $-\Delta_{w,\epsilon}$ in Fourier are

$$\lambda_{k,\epsilon}^2 := \epsilon 4\pi^2 \sum_{j=1}^d \|\xi_j\|^2 + 4\pi^2 \left\| \sum_{j=1}^d \xi_j \right\|^2 \geq \epsilon 4\pi^2 \sum_{j=1}^d \|\xi_j\|^2$$



Special functions

- ▶ Given $\xi := (\xi_1, \dots, \xi_k) \in \mathbb{R}^{kd} := \mathbb{R}^d \times \dots \times \mathbb{R}^d$, set

$$F_\xi^k[m] := \int_{\mathbb{R}^d \times \dots \times \mathbb{R}^d} \exp(-2\pi i \langle \xi, x \rangle) m(dx_1) \cdots m(dx_k)$$

Special functions

- ▶ Given $\xi := (\xi_1, \dots, \xi_k) \in \mathbb{R}^{kd} := \mathbb{R}^d \times \dots \times \mathbb{R}^d$, set

$$F_\xi^k[m] := \int_{\mathbb{R}^d \times \dots \times \mathbb{R}^d} \exp\left(-2\pi i \langle \xi, x \rangle\right) m(dx_1) \cdots m(dx_k)$$

- ▶ $F_\xi^k[\cdot]$ are eigenfunctions of Δ_w

Special functions

- ▶ Given $\xi := (\xi_1, \dots, \xi_k) \in \mathbb{R}^{kd} := \mathbb{R}^d \times \dots \times \mathbb{R}^d$, set

$$F_\xi^k[m] := \int_{\mathbb{R}^d \times \dots \times \mathbb{R}^d} \exp\left(-2\pi i \langle \xi, x \rangle\right) m(dx_1) \cdots m(dx_k)$$

- ▶ $F_\xi^k[\cdot]$ are eigenfunctions of Δ_w
- ▶ The superpositions of the $F_\xi^k[\cdot]$ yield Sobolev functions

Special functions

- ▶ Given $\xi := (\xi_1, \dots, \xi_k) \in \mathbb{R}^{kd} := \mathbb{R}^d \times \dots \times \mathbb{R}^d$, set

$$F_\xi^k[m] := \int_{\mathbb{R}^d \times \dots \times \mathbb{R}^d} \exp\left(-2\pi i \langle \xi, x \rangle\right) m(dx_1) \cdots m(dx_k)$$

- ▶ $F_\xi^k[\cdot]$ are eigenfunctions of Δ_w
- ▶ The superpositions of the $F_\xi^k[\cdot]$ yield Sobolev functions



$$F[m] := \sum_{k=1}^{\infty} \frac{1}{k!} \int_{\mathbb{R}^d \times \dots \times \mathbb{R}^d} a_k(\xi) F_\xi^k[m] d\xi$$

on

$$\mathcal{M} := \mathcal{P}_2(\mathbb{R}^d)$$

when appropriate conditions are imposed on $(a_k)_k$.

- ▶ Say $F \in H^s(\mathcal{M})$ if

$$\sum_{k=1}^{\infty} \frac{1}{k!} \int_{\mathbb{R}^{dk}} |a_k(\xi)|^2 (1 + \lambda_k^2(\xi))^s d\xi < \infty.$$

- ▶ Say $F \in H^s(\mathcal{M})$ if

$$\sum_{k=1}^{\infty} \frac{1}{k!} \int_{\mathbb{R}^{dk}} |a_k(\xi)|^2 (1 + \lambda_k^2(\xi))^s d\xi < \infty.$$

- ▶ The series converges uniformly if there are $C, \delta > 0$ such that

$$|a_k(\xi)| \leq \frac{Ck!}{k^\delta}$$

- ▶ Say $F \in H^s(\mathcal{M})$ if

$$\sum_{k=1}^{\infty} \frac{1}{k!} \int_{\mathbb{R}^{dk}} |a_k(\xi)|^2 (1 + \lambda_k^2(\xi))^s d\xi < \infty.$$

- ▶ The series converges uniformly if there are $C, \delta > 0$ such that

$$|a_k(\xi)| \leq \frac{Ck!}{k^\delta}$$

- ▶ Theorem: If $F \in H^s(\mathcal{M})$ setting

$$U(t, m) := \mathbb{E}\left(F(\mathbb{B}_t^m)\right)$$

- ▶ Say $F \in H^s(\mathcal{M})$ if

$$\sum_{k=1}^{\infty} \frac{1}{k!} \int_{\mathbb{R}^{dk}} |a_k(\xi)|^2 (1 + \lambda_k^2(\xi))^s d\xi < \infty.$$

- ▶ The series converges uniformly if there are $C, \delta > 0$ such that

$$|a_k(\xi)| \leq \frac{Ck!}{k^\delta}$$

- ▶ Theorem: If $F \in H^s(\mathcal{M})$ setting

$$U(t, m) := \mathbb{E}\left(F(\mathbb{B}_t^m)\right)$$

- ▶ we have

$$\partial_t U = \Delta_w U, \quad U(0, \cdot) = F.$$

- ▶ Say $F \in H^s(\mathcal{M})$ if

$$\sum_{k=1}^{\infty} \frac{1}{k!} \int_{\mathbb{R}^{dk}} |a_k(\xi)|^2 (1 + \lambda_k^2(\xi))^s d\xi < \infty.$$

- ▶ The series converges uniformly if there are $C, \delta > 0$ such that

$$|a_k(\xi)| \leq \frac{Ck!}{k^\delta}$$

- ▶ Theorem: If $F \in H^s(\mathcal{M})$ setting

$$U(t, m) := \mathbb{E}\left(F(\mathbb{B}_t^m)\right)$$

- ▶ we have

$$\partial_t U = \Delta_w U, \quad U(0, \cdot) = F.$$

- ▶ For any $r \geq s$, $U(t, \cdot) \in H^r(\mathcal{M})$

Brief review on conservation laws

- ▶ Let us go back to the finite dimensional setting.

Brief review on conservation laws

- ▶ Let us go back to the finite dimensional setting.
- ▶ Let

$H \in C^3(\mathbb{R}^{2d})$ be such that $H(q, \cdot)$ is strictly convex

Brief review on conservation laws

- ▶ Let us go back to the finite dimensional setting.
- ▶ Let

$H \in C^3(\mathbb{R}^{2d})$ be such that $H(q, \cdot)$ is strictly convex

- ▶ Consider the Hamilton–Jacobi equation

$$\partial_t A + H(q, \nabla A) = 0 \quad \text{on} \quad (0, \infty) \times \mathbb{R}^d \times \cdots \times \mathbb{R}^d$$

Brief review on conservation laws

- ▶ Let us go back to the finite dimensional setting.
- ▶ Let

$H \in C^3(\mathbb{R}^{2d})$ be such that $H(q, \cdot)$ is strictly convex

- ▶ Consider the Hamilton–Jacobi equation

$$\partial_t A + H(q, \nabla A) = 0 \quad \text{on} \quad (0, \infty) \times \mathbb{R}^d \times \cdots \times \mathbb{R}^d$$

- ▶ Set $\mathbf{a} := \nabla A$. It is well-known that by differentiation one obtains the conservation laws (when $\epsilon = 0$)

$$\partial_t \mathbf{a} + \nabla \left(H(q, \mathbf{a}) \right) = \epsilon \Delta \mathbf{a} \quad \text{on} \quad (0, \infty) \times \mathbb{R}^d \times \cdots \times \mathbb{R}^d$$

Brief review on conservation laws

- ▶ Let us go back to the finite dimensional setting.
- ▶ Let

$H \in C^3(\mathbb{R}^{2d})$ be such that $H(q, \cdot)$ is strictly convex

- ▶ Consider the Hamilton–Jacobi equation

$$\partial_t A + H(q, \nabla A) = 0 \quad \text{on} \quad (0, \infty) \times \mathbb{R}^d \times \dots \times \mathbb{R}^d$$

- ▶ Set $\mathbf{a} := \nabla A$. It is well-known that by differentiation one obtains the conservation laws (when $\epsilon = 0$)

$$\partial_t \mathbf{a} + \nabla \left(H(q, \mathbf{a}) \right) = \epsilon \Delta \mathbf{a} \quad \text{on} \quad (0, \infty) \times \mathbb{R}^d \times \dots \times \mathbb{R}^d$$

- ▶ Similarly,

$$\partial_t A + H(q, \nabla A) = \epsilon \Delta A \quad \implies \quad \partial_t \mathbf{a} + \nabla \left(H(q, \mathbf{a}) \right) = \epsilon \Delta \mathbf{a}$$

Mean Field Game $n = \infty$

- ▶ We would like to write the analogue of this reasoning in infinite dimension

Mean Field Game $n = \infty$

- ▶ We would like to write the analogue of this reasoning in infinite dimension
- ▶ We define the Hamiltonian

$$\mathcal{H}(m, \zeta) := \int_{\mathbb{R}^d} H(q, \zeta(q)) m(dq)$$

Mean Field Game $n = \infty$

- ▶ We would like to write the analogue of this reasoning in infinite dimension
- ▶ We define the Hamiltonian

$$\mathcal{H}(m, \zeta) := \int_{\mathbb{R}^d} H(q, \zeta(q)) m(dq)$$

- ▶ For a system of n particles

$$m = \frac{1}{n} \sum_{i=1}^n \delta_{x_i} \quad \text{represents the position of particles}$$

Mean Field Game $n = \infty$

- ▶ We would like to write the analogue of this reasoning in infinite dimension
- ▶ We define the Hamiltonian

$$\mathcal{H}(m, \zeta) := \int_{\mathbb{R}^d} H(q, \zeta(q)) m(dq)$$

- ▶ For a system of n particles

$$m = \frac{1}{n} \sum_{i=1}^n \delta_{x_i} \quad \text{represents the position of particles}$$

- ▶ The velocity is

$$\mathbf{v} = \sum_{i=1}^n \xi_i \delta_{x_i}$$

Mean Field Game $n = \infty$

- ▶ We would like to write the analogue of this reasoning in infinite dimension
- ▶ We define the Hamiltonian

$$\mathcal{H}(m, \zeta) := \int_{\mathbb{R}^d} H(q, \zeta(q)) m(dq)$$

- ▶ For a system of n particles

$$m = \frac{1}{n} \sum_{i=1}^n \delta_{x_i} \quad \text{represents the position of particles}$$

- ▶ The velocity is

$$\mathbf{v} = \sum_{i=1}^n \xi_i \delta_{x_i}$$

- ▶ The Hamiltonian restricted to a finite system is

$$\mathcal{H}(m, \mathbf{v}) = \frac{1}{n} H(x_i, \xi_i)$$

- ▶ Consider the Hamilton–Jacobi equation: Find u

$$U^n : (0, T) \times \mathbb{R}^d \times \cdots \times \mathbb{R}^d \rightarrow \mathbb{R}$$

such that

$$\partial_t U^n(t, x) + \mathcal{H}(m^x, \nabla_w U^n(t, x)) = \epsilon_1 \Delta_n U^n(t, x) + \epsilon_2 O_n(U^n(t, x))$$

- ▶ Consider the Hamilton–Jacobi equation: Find u

$$U^n : (0, T) \times \mathbb{R}^d \times \cdots \times \mathbb{R}^d \rightarrow \mathbb{R}$$

such that

$$\partial_t U^n(t, x) + \mathcal{H}(m^x, \nabla_w U^n(t, x)) = \epsilon_1 \Delta_n U^n(t, x) + \epsilon_2 O_n(U^n(t, x))$$

- ▶ Here,

$$\Delta_n = \sum_j^n \Delta_{x_j} + \sum_{j \neq k}^n \text{Trace} \left(\nabla_{x_j x_k} \right), \quad O_n := \sum_j^n \Delta_{x_j}$$

- ▶ Consider the Hamilton–Jacobi equation: Find u

$$U^n : (0, T) \times \mathbb{R}^d \times \cdots \times \mathbb{R}^d \rightarrow \mathbb{R}$$

such that

$$\partial_t U^n(t, x) + \mathcal{H}(m^x, \nabla_w U^n(t, x)) = \epsilon_1 \Delta_n U^n(t, x) + \epsilon_2 O_n(U^n(t, x))$$

- ▶ Here,

$$\Delta_n = \sum_j^n \Delta_{x_j} + \sum_{j \neq k}^n \text{Trace} \left(\nabla_{x_j x_k} \right), \quad O_n := \sum_j^n \Delta_{x_j}$$

- ▶ Let n tend to ∞ to obtain an infinite dimensional equation

Mean Field Game $n = \infty$

- ▶ Coming back to the infinite dimension

Mean Field Game $n = \infty$

- ▶ Coming back to the infinite dimension
- ▶ Recall the Hamiltonian

$$\mathcal{H}(m, \zeta) := \int_{\mathbb{R}^d} H(q, \zeta(q)) m(dq)$$

Mean Field Game $n = \infty$

- ▶ Coming back to the infinite dimension
- ▶ Recall the Hamiltonian

$$\mathcal{H}(m, \zeta) := \int_{\mathbb{R}^d} H(q, \zeta(q)) m(dq)$$

- ▶ Consider on $\mathcal{P}_2(\mathbb{R}^d)$, the Hamilton–Jacobi equation:

$$\partial_t \mathcal{U} + \mathcal{H}(m, \nabla_w \mathcal{U}) = \epsilon_1 \Delta_w \mathcal{U} + \epsilon_2 O(\mathcal{U}) \quad \text{on} \quad (0, T) \times \mathcal{P}_2(\mathbb{R}^d) \quad (2)$$

Mean Field Game $n = \infty$

- ▶ Coming back to the infinite dimension
- ▶ Recall the Hamiltonian

$$\mathcal{H}(m, \zeta) := \int_{\mathbb{R}^d} H(q, \zeta(q)) m(dq)$$

- ▶ Consider on $\mathcal{P}_2(\mathbb{R}^d)$, the Hamilton–Jacobi equation:

$$\partial_t \mathcal{U} + \mathcal{H}(m, \nabla_w \mathcal{U}) = \epsilon_1 \Delta_w \mathcal{U} + \epsilon_2 \mathcal{O}(\mathcal{U}) \quad \text{on } (0, T) \times \mathcal{P}_2(\mathbb{R}^d) \quad (2)$$

- ▶ Set $u := \delta_{L^2} \mathcal{U}$ and apply δ_{L^2} to (2)

Mean Field Game $n = \infty$

- ▶ Coming back to the infinite dimension
- ▶ Recall the Hamiltonian

$$\mathcal{H}(m, \zeta) := \int_{\mathbb{R}^d} H(q, \zeta(q)) m(dq)$$

- ▶ Consider on $\mathcal{P}_2(\mathbb{R}^d)$, the Hamilton–Jacobi equation:

$$\partial_t \mathcal{U} + \mathcal{H}(m, \nabla_w \mathcal{U}) = \epsilon_1 \Delta_w \mathcal{U} + \epsilon_2 O(\mathcal{U}) \quad \text{on } (0, T) \times \mathcal{P}_2(\mathbb{R}^d) \quad (2)$$

- ▶ Set $u := \delta_{L^2} \mathcal{U}$ and apply δ_{L^2} to (2)
- ▶ We obtain the non-local equation

$$\begin{aligned} & \partial_t u + H(m, \nabla u) + \int_{\mathbb{R}^d} \langle \nabla_w u; D_p H(\cdot, \nabla u) \rangle m(dy) = \\ & \epsilon_1 \left(\underline{\Delta_w u} + \Delta u + 2 \int_{\mathbb{R}^d} \operatorname{div} \left[\nabla_w (u[q, m])(y) \right] m(dy) \right) + \epsilon_2 \left(\underline{O(u)} + \Delta u \right) \end{aligned}$$

Mean Field Game $n = \infty$

- ▶ This is an equation introduced in Mean Field Games by P.L. Lions

Mean Field Game $n = \infty$

- ▶ This is an equation introduced in Mean Field Games by P.L. Lions
- ▶ Summary on latest pages:

Mean Field Game $n = \infty$

- ▶ This is an equation introduced in Mean Field Games by P.L. Lions
- ▶ Summary on latest pages:
- ▶ by differentiating an infinite dimensional Hamilton–Jacobi equation, an equation of local nature, one obtains an infinite dimensional perturbed conservation law system

Mean Field Game $n = \infty$

- ▶ This is an equation introduced in Mean Field Games by P.L. Lions
- ▶ Summary on latest pages:
- ▶ by differentiating an infinite dimensional Hamilton–Jacobi equation, an equation of local nature, one obtains an infinite dimensional perturbed conservation law system
- ▶ A side lesson we learnt:

Mean Field Game $n = \infty$

- ▶ This is an equation introduced in Mean Field Games by P.L. Lions
- ▶ Summary on latest pages:
- ▶ by differentiating an infinite dimensional Hamilton–Jacobi equation, an equation of local nature, one obtains an infinite dimensional perturbed conservation law system
- ▶ A side lesson we learnt:
- ▶ the difficulties in conservation laws on \mathbb{R}^d will not go awayon $\mathcal{P}_2(\mathbb{R}^d)$

Mean Field Game $n = \infty$

- ▶ This is an equation introduced in Mean Field Games by P.L. Lions
- ▶ Summary on latest pages:
- ▶ by differentiating an infinite dimensional Hamilton–Jacobi equation, an equation of local nature, one obtains an infinite dimensional perturbed conservation law system
- ▶ A side lesson we learnt:
- ▶ the difficulties in conservation laws on \mathbb{R}^d will not go away on $\mathcal{P}_2(\mathbb{R}^d)$
- ▶

CONCLUSIONS

- ▶ I hope I have convinced you:
 - We have the right concept of Wasserstein Laplacian
 - the Wasserstein Laplacian is important and applicable

CONCLUSIONS

- ▶ I hope I have convinced you:
 - We have the right concept of Wasserstein Laplacian
 - the Wasserstein Laplacian is important and applicable

- ▶ COMPLETES THIS TALK

CONCLUSIONS

- ▶ I hope I have convinced you:
 - We have the right concept of Wasserstein Laplacian
 - the Wasserstein Laplacian is important and applicable

- ▶ COMPLETES THIS TALK

- ▶ Thanks for your attention

CONCLUSIONS

- ▶ I hope I have convinced you:
 - We have the right concept of Wasserstein Laplacian
 - the Wasserstein Laplacian is important and applicable

- ▶ COMPLETES THIS TALK
- ▶ Thanks for your attention
- ▶

CONCLUSIONS

- ▶ I hope I have convinced you:
 - We have the right concept of Wasserstein Laplacian
 - the Wasserstein Laplacian is important and applicable

- ▶ COMPLETES THIS TALK
- ▶ Thanks for your attention
- ▶
- ▶ Outstanding open problem:

CONCLUSIONS

- ▶ I hope I have convinced you:
 - We have the right concept of Wasserstein Laplacian
 - the Wasserstein Laplacian is important and applicable

- ▶ COMPLETES THIS TALK
- ▶ Thanks for your attention
- ▶
- ▶ Outstanding open problem:
- ▶ find the optimal cost for transporting

CONCLUSIONS

- ▶ I hope I have convinced you:
 - We have the right concept of Wasserstein Laplacian
 - the Wasserstein Laplacian is important and applicable
- ▶ COMPLETES THIS TALK
- ▶ Thanks for your attention
- ▶
- ▶ Outstanding open problem:
 - ▶ find the optimal cost for transporting
W GAN from deep learning to
- ▶

CONCLUSIONS

- ▶ I hope I have convinced you:
 - We have the right concept of Wasserstein Laplacian
 - the Wasserstein Laplacian is important and applicable
- ▶ COMPLETES THIS TALK
- ▶ Thanks for your attention
- ▶
- ▶ Outstanding open problem:
- ▶ find the optimal cost for transporting
W GAN from deep learning to
- ▶
- ▶ WILFRID GANGBO in OT
- ▶